## \*HIGHSTABILITY FREQUENCY STANDARDS B ASED ON 199 Hg+IONS 1 N A LINEAR ION TRAP\*

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Microwave-Optical double resonance spectroscopy measurements of the ground state  ${}^2S_{1/2}(F^{\pm}0,m_F^{\pm}O)$  to  ${}^2S_{1/2}(F^{\pm}1,m_{\star})$  hyperfine transition of  ${}^{199}Hg$  ions [1] are used to control a local oscillator providing the most stable atomic frequency standard for averaging times longer than 20,000 seconds. Mercury ions are confined in a Linear lon Trap (1.1T) [2] and collisions with a helium buffer gas cool the icms to near room temperature [3]. Atomic state selection is accomplished by optical pumping using 194 nm light from a  ${}^{202}Hg$  lamp and the 40.507347997x GHz atomic transition is interrogated using Ramsey successive oscillatory fields[4]. With interrogation times of 16 sec.ends a stability of 7x10-]4/T  ${}^{1/2}$  is achieved (Fig.1)[5].

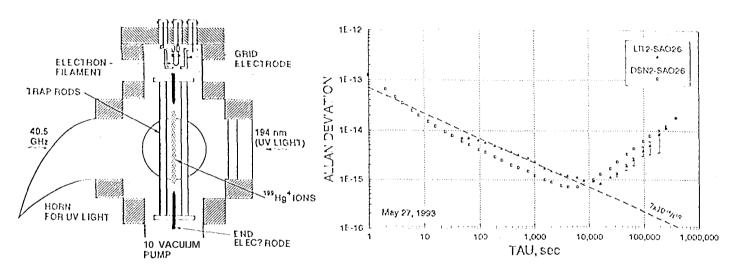


Figure 1: (a) Linear Ion Trap. (b) Ten day stability comparison of LITS-2 and the H-maser SAO-26. The dashed line represents a performance of  $7x10^{-14}/\tau$  lo. Also shown for reference is the stability comparison of two 1 I-masers, SAO-26 and DSN-2 over the same time interval.

Because of the large atomic mass and ground state hyperfine splitting, mercury ions are less susceptible to magnetic and thermal effects than hydrogen or cesium. For good S/N, approximately 107 ions are confined in the trap. Remaining sensitivity to second order Doppler shifts due to thermal and driven motion from the 1f trapping field is the tradeoff for the inherent simplicity in a lamp based, high S/N (i. c large ion cloud), room temperature system. With the addition of a ] 94 nm laser, laser cooling and interrogation on only a few ions would reduce sensitivity to second order Doppler shifts [6].

Stability comparisons are currently underway between two JPL lamp based <sup>199</sup>Hg<sup>4</sup> trapped ion research standards ],] 'J'S-] and LJTS-2, in the present configuration, each standard steers a separate VLG-11 hydrogen maser receiver [7] phase locked to a common hydrogen maser oscillator. In a preliminary 10 day measurement, the short term performance of both standards is o  $_{y}(\tau)$ =1 x10-13/T <sup>1/2</sup>. Figure 2 shows the fractional frequency of the JPL standard LJTS-2 compared against our first research standard LJTS-1.

The deviation from the  $1x 10^{-13}/\tau^{1/2}$  slope between 30,000 and 60,000 is due to limited stability of the control electronics in 1 ITS-1. The long term differential drift between the two Hg+ standards is measured to be less than  $5x10^{\circ}$ ]6/clay. Also shown is a 30 day stability comparison of 1 JTS-2 against the ]]-maser SAO-26 [7]. Beyond 20,000 seconds the measurement is limited by the  $4x10^{-15}$ /day drift of the hydrogen maser.

The [], [] '] 'S-?,/[] [] Comparison is expected to reach 10'16 in 5x10s seconds when electronic upgrading is completed. In the present ion trap standards, the stability floor—is most sensitive to ion number and temperature fluctuations through the second or der Doppler shift [5]. A new extended linear ion trap Configurate ion (LITE) which separates the ion loading and State selection region from the microwave at omic interrogation region is currently under development [8]. This configuration should eliminate most of the Jemaining sensitivity to ion number fluctuations

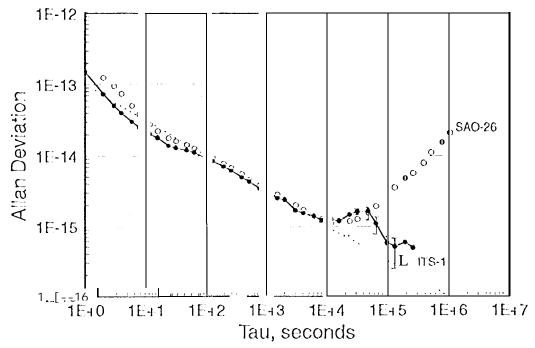


Figure 2: The flat.tional frequency stability of the Hg+ion trap standard I,1'1'S-2 compared against 1.1'1"S-1, and hydrogen maw SAO-26 [6].

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